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CHARACTERIZATION OF HYBRID FERROELECTRIC/HTS THIN FILMS FOR TUNABLE MICROWAVE COMPONENTS

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ABSTRACT

Since the discovery of High-Temperature-Superconductors (HTS) in 1986, a diversity of HTS-based microwave components has been demonstrated. Because of their low conductor losses, HTS-based components are very attractive for integration into microwave circuits for space communication systems. Recent advancements have made deposition of ferroelectric thin films onto HTS thin films possible. Due to the sensitivity of the ferroelectric's dielectric constant (ϵ_r) to an externally applied electric field (E), ferroelectric/superconducting structures could be used in the fabrication of low loss, tunable microwave components. In this paper, we report on our study of $Ba_{0.5}Sr_{0.5}TiO_3/YBa_2Cu_3O_{7-\delta}$ and $Ba_{0.08}Sr_{0.92}TiO_3/YBa_2Cu_3O_{7-\delta}$ ferroelectric/superconducting thin films on lanthanum aluminate (LaAlO3) substrates. For the (Ba:Sr, 0.50:0.50) epitaxial sample, a ϵ_r of 425 and a loss tangent (tan δ) of 0.040 were measured at 298 K, 1.0 MHz, and zero applied E. For the same sample, a ϵ_r of 360 and tan δ of 0.036 were obtained at 77 K, 1.0 MHz, and zero applied E. Variations in ϵ_r from 180 to 360 were observed over an applied E range of 0V/cm≤E≤5.62x10⁴V/cm with little change in tanδ. However, the range of ϵ_r variation for the polycrystalline (Ba:Sr, 0.08:0.92) sample over $0V/cm \le E \le 4.00x10^4V/cm$ was only 3.6 percent while $tan\delta$ increased markedly. These results indicate that a lack of epitaxy between the ferroelectric and superconducting layers decreases tuning and increases microwave losses.

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INTRODUCTION

Since the discovery of High-Temperature-Superconductors (HTS) in 1986, it has been possible to fabricate HTS-based planar microwave components such as resonators and filters¹. The use of HTS films in place of normal conductors (e.g. gold and copper) has greatly reduced conductor losses and consequently insertion losses². Therefore, significant attention has been given to the quality and processing of HTS films as efforts continue towards the inclusion of HTS-based circuits in working systems. One of the areas that has received considerable attention for the insertion of HTS-based components is that of tunable microwave circuits³. This requires the development of hybrid superconductor/ferroelectric structures. Ferroelectric materials have a dielectric constant (ϵ_{τ}) that can be altered by an externally applied electric field (E). Therefore, optimization of these structures' material properties and careful characterization of their electrical properties should result in high quality, low loss tunable microwave components, such as capacitors, filters, and phase shifters. In this paper, we report on our study of Ba_{0.8}Sr_{0.5}TiO₃/YBa₂Cu₃O_{7.8} and Ba_{0.08}Sr_{0.92}TiO₃/YBa₂Cu₃O_{7.8} ferroelectric/superconducting thin films on lanthanum aluminate (LaAlO₃) substrates.

EXPERIMENTAL

The $Ba_{0.5}Sr_{0.5}TiO_3/YBa_2Cu_3O_{7.6}/LaAlO_3$ (BST/YBCO/LaAlO_3) structure which was investigated consisted of an 800 nm thick $Ba_{0.5}Sr_{0.5}TiO_3$ film, known hereinafter as BST, deposited on a 300 nm thick $YBa_2Cu_3O_{7.6}$ film which coated a 5.0 mm×4.0 mm×0.50 mm LaAlO_3 substrate. Both layers were deposited "in-situ" using laser ablation. To elucidate the effects of film microstructure on tuning and loss, metalorganic deposited $Ba_{0.08}Sr_{0.92}TiO_3$ films, known hereinafter as MOD BST films, were prepared at room temperature and post annealed at $700^{\circ}C$ to produce polycrystalline films⁴. After deposition of the multilayer structures,

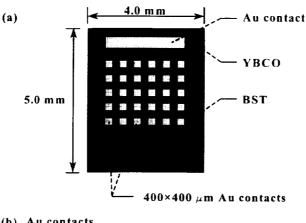
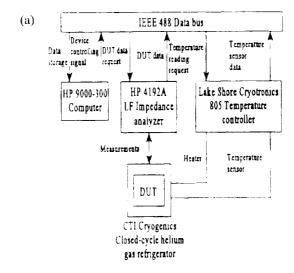




Figure 1. A schematic representation of the Ba_xSr_{1.x}TiO₃(x=0.50,0.08)/YBa₂Cu₃O_{7.8}/LaAlO₃ structure. (a) Top view, (b) Side view.



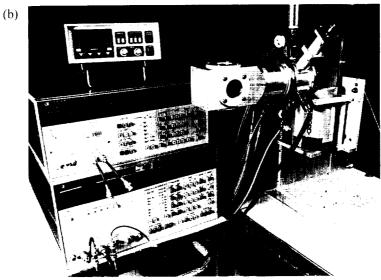


Figure 2. Experimental configuration. (a) Schematic representation, (b) Actual setup.

transition temperatures (T_c) above 89 K were measured for the YBCO films using a four point probe technique indicating that no degradation of the these films occurred as a result of the ferroelectric film deposition. Thirty $400\mu\text{m}\times400\mu\text{m}\times2.5\mu\text{m}$ gold (Au) contacts were fabricated on the ferroelectric layer of the structures using standard photolithography and wet etching techniques. In addition, a Au electrode was deposited onto a portion of the structures' YBCO layer, exposed by etching the ferroelectric layer away with a 7.0 percent solution of hydrofluoric (HF) acid. Figure 1 shows a schematic representation of the aforementioned structures. This configuration was used to measure the capacitance (C) at 1.0 MHz from which the ϵ_r of the ferroelectric layer was thereby calculated using Eq. (1),

$$\epsilon_r \cdot \frac{Cd}{\epsilon_s A}$$
 (1)

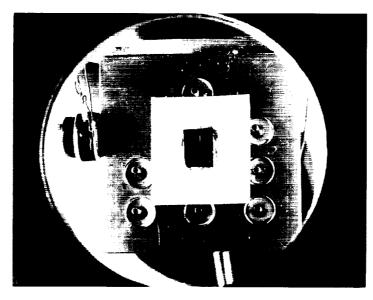


Figure 3. Sample mounted inside vacuum chamber showing electrical feedthroughs.

where A represents the area of the gold contact, d represents the thickness of the ferroelectric layer, and $\epsilon_{\rm o}$ =8.854×10⁻¹² F/m represents the permittivity of free space.

Values of $\epsilon_{\rm r}$ for the ferroelectric layers were obtained as a function of temperature and dc bias. An HP-4192A LF impedance analyzer, a CTI-Cryogenics closed-cycle, helium gas refrigerator, and a Lake Shore Cryotronics 805 temperature controller all linked through an IEEE-488 data bus to an HP 9000-300 computer served as the main components of the experimental configuration, shown in Figure 2, which performed these measurements. As seen in Figure 3, the sample, known as the device-under-test (DUT), was mounted on a sample holder having electrical feedthroughs for dc bias. The temperature sensor, a silicon diode, was bolted directly to the sample holder. The sample holder was bolted to the cold head of the refrigerator to allow for measurements at cryogenic temperatures. An HP-Basic computer program was written to automate the measurement system allowing the user to take data as a function of temperature at a constant dc voltage and as a function of dc voltage at a constant temperature. All measurements were taken at "user-input" pre-selected conditions. Preselected dc voltage conditions were directly outputted to the impedance analyzer and

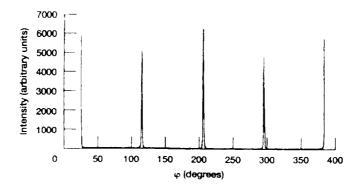


Figure 4. Phi scan about the BST (420) peak. Four-fold symmetry shows good in-plane alignment of the BST film on the YBCO film.

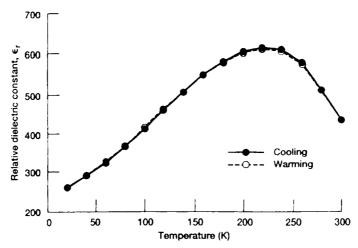


Figure 5. Dielectric constant versus temperature for the BST (Ba:Sr, 0.50:0.50) layer in the BST/YBCO/LaAlO₃ structure at 1.0 MHz and zero E.

temperatures conditions were outputted as "set-points" to the temperature controller which controlled and monitored the temperature of the DUT. When specified conditions were met, data on the DUT's capacitance and loss tangent (tan δ) were measured by the impedance analyzer and stored along with the DUT's $\epsilon_{\rm r}$ calculated by the computer program using Eq. 1, to a 3.5 inch disk for further analysis.

RESULTS

The ferroelectric layers were analyzed using x-ray diffraction (XRD). As seen in Figure 4, Phi scans about the (420) peak of the BST film show in-plane epitaxy with the underlying YBCO film. However, diffraction peaks for the MOD BST film were very weak at high angles and there was no evidence of in-plane epitaxy.

For the electrical characterization of the samples, the first test undertaken focused on measuring $\epsilon_{\rm r}$ as a function of temperature at a frequency of 1.0 MHz and at zero E. Measurements were taken at intervals of 20 K during cooling from 300 to 20 K and during

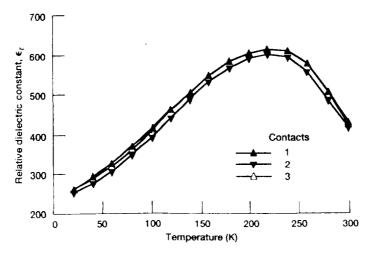


Figure 6. Dielectric constant values of different contacts versus temperature for the BST (Ba:Sr, 0.50:0.50) layer in the BST/YBCO/LaAlO, structure at 1.0 MHz and zero E.

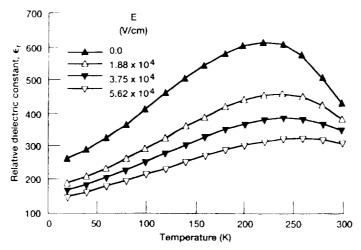


Figure 7. Dielectric constant versus temperature for the BST (Ba:Sr, 0.50:0.50) layer in the BST/YBCO/LuAIO, structure at 1.0 MHz and several values of E.

warming from 20 to 300 K to determine if the samples were sensitive to the temperature cycling. At room temperature, values of $\epsilon_{\rm r}$ and $\tan\delta$ for the BST sample of 425 and 0.040, respectively, were obtained. Neither the $\epsilon_{\rm r}$ data, seen in Figure 5, nor the $\tan\delta$ data showed any signs of temperature hysteresis as a result of the temperature cycling. Figure 6 shows $\epsilon_{\rm r}$ data from different contacts across the sample's surface taken under the same temperature and bias conditions. This test was performed to determine the degree of uniformity of the BST film. Note that the obtained results are quantitatively and qualitatively very similar. This suggests excellent composition and structural (i.e., thickness) uniformity throughout the BST layer. The behavior of $\epsilon_{\rm r}$ and $\tan\delta$ as a function of E was also investigated. The results are shown in Figure 7 and Figure 8. In Figure 7, each curve represents the $\epsilon_{\rm r}$ of the BST layer as a function of temperature at a particular E value. Note that there is a decrease in $\epsilon_{\rm r}$ as the magnitude of the applied field is increased. Maximum variation in $\epsilon_{\rm r}$, from $\epsilon_{\rm r}$ =610 to $\epsilon_{\rm r}$ =300, occurred at a temperature near 220 K. For practical microwave components, such as tunable

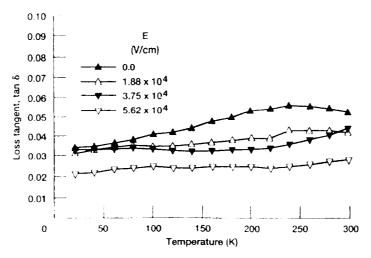


Figure 8. Loss tangent versus temperature for the BST (Ba:Sr, 0.50:0.50) layer in the BST/YBCO/LaAlO₃ structure at 1.0 MHz and several values of **E**.

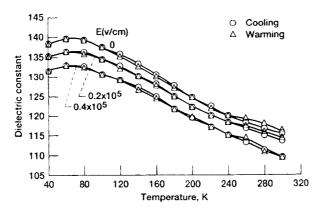


Figure 9. Dielectric constant versus temperature for a polycrystalline, 500 nm thick, metalorganic deposited MOD BST (Ba:Sr, 0.08:0.92) layer in the BST/YBCO/LaAlO₃ structure at 1.0 MHz and several values of E.

capacitors, filters, and phase shifters, a change in ϵ_{τ} from 300 to 610 will create a frequency range within which the component can be tuned⁵⁻⁷. At 77 K, a temperature at which YBCO is in the superconducting state, values of ϵ_{τ} changed from 180 to 360 in the range of $0V/cm \le E \le 5.62 \times 10^4 V/cm$. At this temperature, both the HTS film's low conductor losses as well as the ferroelectric's tunability could be fully exploited.

Figure 8 shows results of $\tan\delta$ as a function of temperature and **E**. It can be seen that the values for $\tan\delta$ vary slightly as a function of **E** in the aforementioned field range. This result is very important for microwave applications since it shows that the ϵ_r can be altered without dramatically enhancing the ferroelectric's microwave losses. The range of tunability for the polycrystalline, 500 nm thick, MOD BST film over $0V/\text{cm} \le E \le 4.0 \times 10^4 V/\text{cm}$ at 60 K was only 3.6 percent, see Figure 9. In addition, $\tan\delta$ was very sensitive to **E** as shown in Figure 10. We dismiss the possibility that the lower tuning in the MOD BST film is due to the composition differences since measurements on epitaxial STO and $Ba_{0.10}Sr_{0.90}TiO_3$ films deposited on YBCO displayed ϵ_r values that could be decreased by 30 and 10 percent respectively under the same electric field. Therefore, the fact that the tuning was lower for the MOD BST film suggests that the lack of epitaxy decreases tuning and also increases microwave losses.

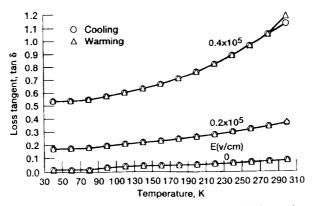


Figure 10. Loss tangent versus temperature for a polycrystalline, 500 nm thick, metalorganic deposited MOD BST (Ba:Sr, 0.08:0.92) layer in the BST/YBCO/LaAlO₃ structure at 1.0 MHz and several values of **E**.

CONCLUSION

A ϵ_r of 425 and a tan δ of 0.040 were measured at 298 K, 1.0 MHz, and zero **E** for an epitaxial BST ferroelectric layer. At 77 K, 1.0 MHz, and zero **E**, a ϵ_r of 360 and a tan δ of 0.036 were measured. Maximum change in ϵ_r , from 300 to 610, was observed at 220 K in the **E** range of 0V/cm \leq **E** \leq 5.62x10⁴V/cm with minimal change in tan δ . The range of tunability for a polycrystalline, 500 nm thick, MOD BST film was very small, 3.6 percent, and tan δ increased markedly with **E** over 0V/cm \leq **E** \leq 4.0x10⁴V/cm. This suggests that the lack of epitaxy could adversely affect the tuning range and the microwave losses of ferroelectric/HTS based microwave components.

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